

## Comment on “Electronic fine structure in the electron-hole plasma of $\text{SrB}_6$ ”

In a recent Letter [1], Rodriguez *et al.* addressed the origin of the high-temperature weak ferromagnetism found in doped hexaborides [2] by calculating details of the band structure of  $\text{SrB}_6$  in the vicinity of the X-point in the Brillouin zone. They conclude that crystal field induced electron-hole mixing has an important effect on the density of states (DOS) and suggest nesting between electron Fermi pockets as an alternative explanation to the excitonic ferromagnetism [3]. We would like to question several conclusions in this work [1].

(i) Rodriguez *et al.* claim that the two bands cross only along the symmetry lines  $(\xi, 0, 0)$  and  $(\frac{\pi}{a}, \xi, 0)$ . This result contradicts a previous band structure study, which found band crossing in *all symmetry planes* passing through the X-point [4]. The latter conclusion has a simple symmetry interpretation. At the X-point the valence and conduction bands belong to different irreducible representations of the small group:  $X_3$ , basis function  $yz$ , and  $X'_3$  symmetry, basis function  $x(y^2 - z^2)$ , respectively. The band mixing occurs away from the X-point and is given by the product  $X_3 \times X'_3$

$$\Delta_{\text{cryst}} \simeq E_b k_x k_y k_z (k_y^2 - k_z^2) / b^5, \quad (1)$$

where  $E_b$  is a bandwidth and  $b = \frac{2\pi}{a}$  is a reciprocal lattice vector. The crystal field hybridization is strongly suppressed by the high power of momentum in Eq. (1) ( $k_F \sim 0.1b$ ). Therefore, not only is the hybridization more anisotropic than suggested in Ref. [1], but also it is very weak on the scale of the band overlap. Our TB-LMTO calculations [3] as well as FLAPW method [5] yield an extremely small band mixing for a general direction in the Brillouin zone in agreement with Eq. (1).

(ii) Even for a strong hybridization, the DOS is quite different from the results of Ref. [1]. Fig. 1 shows the DOS for two parabolic bands with anisotropic effective masses [1,3], band overlap  $E_G = 100$  meV, and two angular dependences for the mixing matrix element: isotropic, which models an excitonic gap [3], and anisotropic given by Eq. (1). The amplitudes were chosen to give the same maximum splitting of  $2\Delta_{\text{max}} = 15$  meV. Because of the anisotropic dispersion the DOS is slightly reduced at the Fermi level. The true gap in the DOS opens only for a sufficiently large excitonic gap  $\Delta > 0.16E_G$ . Note that the tetrahedron method for DOS does not allow for crossing of bands, which are numbered in increasing order of energy at every k-point, and, therefore, may lead to spurious peaks [1] due to artificial anticrossings.

(iii) Correlation effects in the  $e$ - $h$  plasma lead to a forbidden range for the band overlap values in the hexaborides [6]. At the first order metal-insulator transition a semiconductor jumps into a dense plasma with  $r_s \approx 1$ . Rodriguez *et al.* have estimated  $r_s = 2.5$  for the band overlap  $E_G = 100$  meV in  $\text{SrB}_6$ , a value lying within the

forbidden range. The real overlap must be larger. Therefore, the conclusion that 0.5 % doping fills hole Fermi pockets in  $\text{SrB}_6$  [1], which is based on the above underestimated value of  $E_G$ , is by no means reliable and does not exclude electron-hole pairing at this doping level.

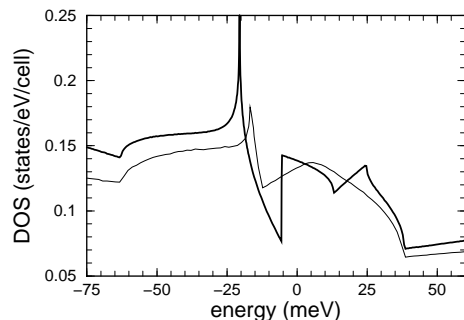


FIG. 1. The density of states in  $\text{SrB}_6$  for two types of interband hybridization (a) isotropic excitonic gap, bold line, and (b) anisotropic crystal-field mixing Eq. (1), thin line.

(iv) Nesting of electron pockets between inequivalent X-points was suggested in Ref. [1] as another source of CDW and SDW instabilities, which does not involve holes from the valence band. This suggestion is physically incorrect. The nesting condition refers to a sign change under translation in momentum space:  $\epsilon(k+Q) \approx -\epsilon(k)$ . Only in such a case can repulsive  $e$ - $e$  interaction produce a density-wave instability in  $e$ - $h$  channel, see e.g. [7].

Lastly, the interpretation of the de Haas-van Alphen measurements quoted in [1] seems to us to be problematic in view of the high values of the low temperature resistivity reported for  $\text{SrB}_6$  [8] and its strong dependence on stoichiometry. Further information is required before a definitive conclusion can be made.

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- [1] C. O. Rodriguez, R. Weht, and W. E. Pickett, Phys. Rev. Lett. **84**, 3903 (2000).
  - [2] D. P. Young *et al.*, Nature **397**, 412 (1999).
  - [3] M. E. Zhitomirsky, T. M. Rice, and V. I. Anisimov, Nature **402**, 251 (1999); cond-mat/9904330.
  - [4] A. Hasegawa and A. Yanase, J. Phys. C **12**, 5431 (1979).
  - [5] S. Massidda *et al.*, unpublished (2000).
  - [6] M. E. Zhitomirsky and T. M. Rice, Phys. Rev. B **62**, No. 3 (2000); cond-mat/9910272.
  - [7] P. Fazekas, *Lecture Notes on Electron Correlation and Magnetism*, World Scientific (1999).
  - [8] H. R. Ott *et al.*, Z. Phys. B **102**, 83 (1997).